

# Focal mechanism determination using high-frequency waveform matching and its application to small magnitude induced earthquakes

Junlun Li, Haijiang Zhang, H. Sadi Kuleli and M. Nafi Toksoz

Earth Resources Laboratory, Department of Earth, Atmospheric and Planetary Sciences, Massachusetts Institute of Technology, 77 Massachusetts Avenue, Cambridge, MA 02139, USA. E-mail: junlunli@mit.edu

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### SUMMARY

We present a new method using high-frequency full waveform information to determine the focal mechanisms of small, local earthquakes monitored by a sparse surface network. During the waveform inversion, we maximize both the phase and amplitude matching between the observed and modelled waveforms. In addition, we use the polarities of the first *P*-wave arrivals and the average S/P amplitude ratios to better constrain the matching. An objective function is constructed to include all four criteria. An optimized grid search method is used to search over all possible ranges of source parameters (strike, dip and rake). To speed up the algorithm, a library of Green's functions is pre-calculated for each of the moment tensor components and possible earthquake locations. Optimizations in filtering and cross correlation are performed to further speed the grid search algorithm. The new method is tested on a five-station surface network used for monitoring induced seismicity at a petroleum field. The synthetic test showed that our method is robust and efficient to determine the focal mechanism when using only the vertical component of seismograms in the frequency range of 3–9 Hz. The application to dozens of induced seismic events showed satisfactory waveform matching between modelled and observed seismograms. The majority of the events have a strike direction parallel with the major NE-SW faults in the region. The normal faulting mechanism is dominant, which suggests the vertical stress is larger than the horizontal stress.

**Key words:** Earthquake source observations; Seismicity and tectonics; Body waves; Wave propagation.

## **1 INTRODUCTION**

In this study we use high-frequency seismograms for determining the focal mechanisms of small earthquakes recorded by a sparse network of seismic stations. The method involves calculating synthetic seismograms for a series of moment tensors and finding the source mechanism where the observed and synthetic seismograms match the best. The method is especially useful for determining the source mechanisms of induced seismic events.

Induced seismicity is a common phenomenon in oil/gas reservoirs and in reservoirs where activities are in progress and internal stress distribution is changed due to water injection, fluid extraction, etc. (Rutledge & Phillips 2003; Rutledge *et al.* 2004; Chan & Zoback 2007; Deichmann & Giardini 2009; Bischoff *et al.* 2010; Segall 2010; Suckale 2010). For example, the gas/oil extraction can cause reservoir compaction and reactivate pre-existing faults and induce earthquakes (e.g. Chan & Zoback 2007; Miyazawa *et al.* 2008; Sarkar *et al.* 2008). By studying the patterns of the induced seismicity over an extended time period (e.g. location and focal mechanism), a time-lapse history of the stress changes in the fields may be reconstructed.

Induced earthquakes usually have small magnitudes and are generally recorded at sparse local stations. As a result, it is difficult to have enough seismic waveforms with high signalto-noise ratio for picking the polarity information of first *P*wave arrivals. Therefore, it is challenging, if not impossible, to use only the *P*-wave polarity information (even when adding *S*/*P* amplitude ratios) as used in conventional methods to constrain the focal mechanisms of the induced earthquakes (e.g. Hardebeck & Shearer 2002, 2003). Waveform matching has been used to determine earthquake focal mechanisms on a regional and global scale using low-frequency waveform information (Zhao & Helmberger 1994; Zhu & Helmberger 1996; Pasyanos *et al.* 1996; Dreger *et al.* 1998; Tan & Helmberger 2007). However, in induced seismicity cases, waveforms usually have higher frequencies.

High-frequency waveform matching, in addition to polarity information, has been used to determine the focal mechanism of induced earthquakes in a mine with a dense network of 20 stations (Julià *et al.* 2009). They used a constant velocity model to calculate the Green's functions, and performed the focal mechanism inversion in frequency domain without phase information in a least-square sense between the synthetic and filtered observed data generally below 10 Hz.

To retrieve reliable solutions, our study uses high-frequency, full waveform information (both *P* and *S*) to determine the focal mechanism of small earthquakes recorded by a sparse five-station network.

Using the known velocity structure, we calculate the Green's functions for all moment tensor components of the source and for each location (hypocentre) and then the synthetic seismograms. To find the best match between the observed and synthetic seismograms, we formulate an objective function that incorporates information from different attributes in the waveforms: the cross-correlation values between the modelled waveforms and the data, the L<sub>2</sub> norms of the waveform differences, the polarities of the first *P* arrivals and the *S/P* average amplitude ratios.

For real application we use data from a five-station network monitoring induced seismicity at a petroleum field (Sarkar 2008; Zhang *et al.* 2009). The method is tested with data from a synthetic seismic event and then applied to real events selected from an oil and gas field.

#### 2 METHOD

The focal mechanism can be represented by a  $3 \times 3$  second-order moment tensor (Stein & Wysession 2003). Usually, as there is no rotation of mass involved in the rupturing process, the tensor is symmetric and has only six independent components. Here we assume the focal mechanism of the small induced events can be represented by pure double couples (DCs, Rutledge & Phillips 2002), though it is possible that a volume change or Compensated Linear Vector Dipoles (CLVD) part may also exist. The constraining of focal mechanism as DC has additional advantages. For example, anisotropy, which is not included by the isotropic Green's functions, often raises spurious non-DC components in focal mechanism determination. By assuming focal mechanism to be DC, this spurious non-DC component can be ruled out while the DCs can be well recovered (Šílený & Vavryčuk 2002). In our analysis we describe the source in terms of its strike, dip and rake, and determine DC components from these three parameters. For each component of a moment tensor, we can use the Discrete Wavenumber Method (DWN) (Bouchon 1981, 2003) to calculate its Green's functions  $G_{iik}^{n}(t)$ . Using reciprocity by strain Green's tensors can further improve the efficiency of calculating the Green's functions, especially when the sources greatly outnumber the stations (Zhao et al. 2006). The structure between the earthquake and the station is represented as a 1-D layered medium. The modelled waveform from a certain combination of strike, dip and rake is expressed as a linear combination of weighted Green's functions:

$$V_i^n = \sum_{j=1}^3 \sum_{k=1}^3 m_{jk} G_{ij,k}^n(t) * s(t),$$
(1)

where  $V_i^n$  is the modelled *i*-th (north, east or vertical) component at station *n*;  $m_{jk}$  is the moment tensor component and is determined by the data from all stations;  $G_{ij,k}^n(t)$  is the *i*-th component of the Green's functions for the (j, k) entry at station *n*, and s(t) is the source time function. In this study, a smooth ramp is used for s(t), whose duration can be estimated from the spectra of the recorded seismograms (Bouchon 1981). In the synthetic and real data cases, the duration of s(t) is about 0.1 s.

Earthquake locations are usually provided by the traveltime location method. However, due to uncertainties in velocity model and arrival times, the seismic event locations may have errors, especially in focal depth. While matching the modelled and observed waveforms, we also search for an improved location around the catalogue location.

Before the grid search is performed, we build a Green's functions library. We pre-calculate all  $G_{ii,k}^n(t)$ 's for all possible event locations

at all stations and store them on the disk. When we perform the grid search, we simply need to do a linear combination of  $G_{ij,k}^n(t) * s(t)$ , each of which is weighted by  $m_{ik}$ .

To determine the best solution, we construct an objective function that characterizes the similarity between the modelled and observed waveforms. We use the following objective function, which evaluates four different aspects of the waveform information:

$$\begin{aligned} & \underset{n=1}{\operatorname{maximize}} \left[ J(x, y, z, \operatorname{str}, \operatorname{dip}, \operatorname{rake}, ts) \right] = \\ & \sum_{n=1}^{N} \sum_{j=1}^{3} \left\{ \alpha_{1} \max \left( \tilde{d}_{j}^{n} \otimes \tilde{v}_{j}^{n} \right) - \alpha_{2} \left\| \tilde{d}_{j}^{n} - \tilde{v}_{j}^{n} \right\|_{2} \right. \\ & + \alpha_{3} f \left[ \operatorname{pol}(\tilde{d}_{j}^{n}), \operatorname{pol}(\tilde{v}_{j}^{n}) \right] + \alpha_{4} h \left[ \operatorname{rat} \left( \frac{S(d_{j}^{n})}{P(d_{j}^{n})} \right), \operatorname{rat} \left( \frac{S(v_{j}^{n})}{P(v_{j}^{n})} \right) \right] \right\} . \end{aligned}$$

Here  $\tilde{d}_j^n$  is the normalized data and  $\tilde{v}_j^n$  is the normalized modelled waveform. Since it is difficult to obtain accurate absolute amplitudes due to site effects in many situations, we normalize the filtered observed and modelled waveforms before comparison. The normalization used here is the energy normalization, such that the energy of the normalized wave train within a time window adds to unity. Compared to peak amplitude normalization, energy normalization is less affected by site effect, which may cause abnormally large peaks due to focusing and other factors. In a concise form, this normalization can be written as

$$\tilde{d}_j^n = \frac{d_j^n}{\sqrt{\int_{t_1}^{t_2} \left(d_j^n\right)^2 \mathrm{d}t}},\tag{3}$$

where  $t_1$  and  $t_2$  are the boundaries of the time window.

The objective function J in eq. (2) consists of four terms.  $\alpha_1 - \alpha_4$ are the weights for each term. Each weight is a positive scalar number and is optimally chosen in a way such that no single term will overdominate the objective function. The first term in eq. (2) evaluates the maximum cross correlation between the normalized data  $(\tilde{d}_i^n)$  and the normalized modelled waveforms  $(\tilde{v}_i^n)$ , which is mostly sensitive to phase differences. From the cross correlation, we find the time-shift to align the modelled waveform with the observed waveform. In high-frequency waveform comparisons, cycle-skip is a special issue requiring extra attention: overshifting the waveform makes the wiggles in the data misalign with wiggles of the next cycle in the modelled waveforms. Therefore, allowed maximum time-shift should be pre-determined by the central frequency of the waveforms. The second term evaluates the L<sub>2</sub> norm of the direct differences between the aligned modelled and observed waveforms (note the minus sign of the second term to minimize the amplitude differences, which measures both phase and amplitude differences). The reason for maximizing the cross-correlation value and minimizing the direct difference between the observed and modelled waveforms is to match both the similarity of the waveforms and the actual amplitudes. The first two terms are not independent to each other, however, they have different sensitivities at different frequency bands and by combining them together the waveform similarity can be better characterized. The third term evaluates whether the polarities of the first P-wave arrivals as observed in the data are consistent with those in the modelled waveforms. 'pol' is a weighted sign function which can be  $\{\beta, -\beta, 0\}$ , where  $\beta$  is a weight reflecting our confidence in picking the polarities of the first P-wave arrivals in the observed data. Zero (0) means undetermined polarity. f is a function that penalizes the polarity sign inconsistency in

such a way that the polarity consistency gives a positive value while polarity inconsistency gives a negative value. The matching of the first *P*-wave polarities between modelled and observed waveforms is an important condition for determining the focal mechanism, when the polarities can be clearly identified. Polarity consistence at some stations can be violated if the polarity is not confidently identified (small  $\beta$ ) and the other three terms favour a certain focal mechanism. Therefore, the polarity information is incorporated into our objective function with some flexibility. By summing over the waveforms in a narrow window around that arrival time and checking the sign of the summation, we determine the polarities robustly for the modelled data. For the observed data, we determine the *P*-wave polarities manually.

The S/P amplitude ratio is also very important in determining the focal mechanism. The fourth term in the objective function is to evaluate the consistency of the S/P amplitude ratios in the observed and modelled waveforms (Hardebeck & Shearer 2003). The 'rat' is the ratio evaluation function and it can be written as

$$\operatorname{rat} = \frac{\int_{T_2}^{T_3} |r_j^n(t)| \, \mathrm{d}t}{\int_{T_2}^{T_2} |r_j^n(t)| \, \mathrm{d}t},\tag{4}$$

where  $[T_1 T_2]$  and  $[T_2 T_3]$  define the time window of *P* and *S* waves, respectively, and  $r_j^n$  denotes either  $d_j^n$  or  $v_j^n$ . The term *h* is a function, which penalizes the ratio differences so that the better matching gives a higher value. Note that here we use the un-normalized waveforms  $d_j^n$  and  $v_j^n$ .

In general, the amplitudes of P waves are much smaller than those of S waves. To balance the contribution between P and S waves, we need to fit P and S waves separately using the first two terms in eq. (2). Also, by separating S from P waves and allowing an independent time-shift in comparing observed data with modelled waveforms, it is helpful to deal with incorrect phase arrival time due to incorrect Vp/Vs ratios (Zhu & Helmberger 1996). We calculate both the first P and S arrival times by the finite difference eikonal solver (Podvin & Lecomte 1991). The wave train is then separated into two parts at the beginning of the S wave. To reduce the effect of uncertainty in the origin time, we first align the modelled and observed data using first arrivals and then define the alignment by cross correlation.

The processing steps can be summarized as follows:

(1) use the known velocity structure to generate a Green's function library;

(2) calculate the first *P*- and *S*-wave arrival times using the finitedifference traveltime solver (Podvin & Lecomte 1991), and separate the *S*-wave segment from the *P*-wave segment according to the traveltime information;

(3) for separated *P*- and *S*-wave segments:

(i) determine the time-shift by cross correlation between modelled and observed data,

(ii) evaluate the maximum cross-correlation value and  $L_2$  norm between the aligned modelled and observed data,

(iii) identify the first arrival polarities and

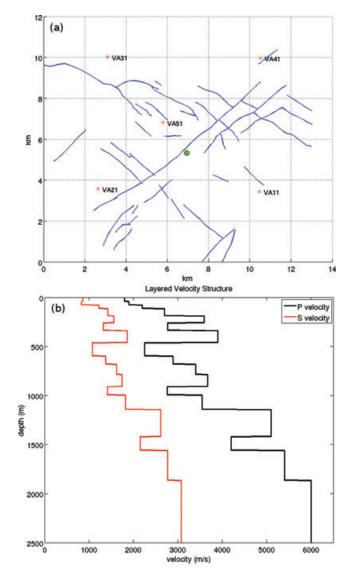
(iv) calculate the average amplitudes of the P- and S-wave segments, and

(4) determine the best fit mechanism by maximizing the objective function.

To find the similarity between modelled and observed waveforms, we do two kinds of basic computations: filtering and cross correlation. These two computations are very time consuming when millions of modelled traces are processed. To expedite the computation we use the following manipulations:

$$v_i^n = F * V_i^n = F * \sum_{j=1}^3 \sum_{k=1}^3 m_{jk} G_{ij,k}^n(t) * s(t)$$
  
=  $\sum_{j=1}^3 \sum_{k=1}^3 m_{jk}^n [F * (G_{ij,k}^n(t) * s(t))],$  (5)

$$d_{i} \otimes v_{i} = d_{i} \otimes \sum_{j=1}^{3} \sum_{k=1}^{3} m_{jk} F * G_{ij,k} * s(t)$$
  
=  $\sum_{j=1}^{3} \sum_{k=1}^{3} m_{jk} [d_{i} \otimes (F * G_{ij,k} * s(t))],$  (6)



**Figure 1.** (a) Locations of stations (red stars) and the synthetic event (green dot). The identified faults, stations and the synthetic event are plotted in a local reference system. (b) Layered velocity structure for *P* and *S* waves for this region. Velocity data are derived from nearby well logs.

where *F* denotes the impulse response of a filter; '\*' denotes time domain convolution;  $d_i^n$  and  $v_i^n$  denote the *i*-th component of the filtered observed and modelled data at station *n*, respectively; ' $\otimes$ ' denotes the cross correlation. These two equations indicate that we can apply the filtering and cross correlation into the summation to avoid filtering and cross correlation repetitively during the search over all strikes, dips and rakes. A large amount of time is thereby saved, and the searching speed is boosted by an order of magnitude.

By pre-calculating the library of Green's functions and manipulating the filtering and cross correlation, we greatly speed up the grid search process. Searching through all possible X, Y and Z for location and strikes, dips and rakes for focal mechanisms often results in over 10 million different waveforms to be compared with the data. Since the grid search can be easily parallelized, it can be done on a multicore desktop machine within 10 min. The computation of the Green's function library using DWN takes more time, but it only needs to be computed once.

### **3 SYNTHETIC DATA TEST**

We first test the accuracy and robustness of our method using a synthetic data set. We use the same station distribution and velocity model from a field shown in Fig. 1 (Sarkar 2008). The layered velocity model is obtained from nearby well logs.

We use the DWN method to generate clean synthetic seismograms and manually pick the first *P* arrival times, as we will do on the real field data. A source located at 1227 m beneath the surface, with a focal mechanism of strike of  $210^{\circ}$ , dip of  $50^{\circ}$  and rake of  $-40^{\circ}$ , is used to generate the synthetic data. The synthetic seismograms are shown in Fig. 2. We use a horizontal grid spacing of 150 m and vertical grid spacing of 50 m with the search range

**Table 1.** Statistics of focal mechanism parameters in the 'clean' synthetic test. The hypocentre is recoordinated as (0, 0, 0) m for discussing the mean and standard deviation, and the source has a strike of  $210^{\circ}$ , dip of  $50^{\circ}$  and rake of  $-40^{\circ}$ . The best solution has the identical focal parameters with the source.

	Strike (°)	Dip (°)	Rake (°)	<i>X</i> (m)	<i>Y</i> (m)	<i>Z</i> (m)
Mean	203	47	-45	-30	10	0
Std.	13	9	11	110	130	0

of -900 m  $\leq X,Y \leq$  900 m and -400 m  $\leq Z \leq$  400 m. The reason for choosing smaller vertical grid spacing is that the seismograms are very sensitive to focal depth. Any vertical shift in hypocentre changes the multiple reflection and refraction patterns. The frequency band used in our study is 3-9 Hz, which contains the dominant energy in waveforms recorded from typical induced earthquakes in that field. The searching interval in strike, dip and rake is 10° in this test and, hereafter, in the real data test. This spacing choice indicates that our resolution is 5° at best. Because the auxiliary plane solution and the fault plane solution give the identical waveform, this means that half of the model space  $[0^{\circ} \leq$ strike  $\leq 360^\circ$ ;  $0^\circ \leq dip \leq 90^\circ$ ;  $-180^\circ \leq rake \leq 180^\circ$ ] is redundant. Therefore, by constraining the model space in  $[0^\circ \le \text{strike} \le 360^\circ]$ ;  $0^{\circ} \leq dip \leq 90^{\circ}; -90^{\circ} \leq rake \leq 90^{\circ}]$  (Zhao & Helmberger 1994), we can eliminate the redundancy and further shorten the search time by half. The weights  $\alpha_1 - \alpha_4$  in the objective function (eq. 2) were tried with different values, and we selected ones that balance different terms. We used  $\alpha_1 = 3$ ,  $\alpha_2 = 3$ ,  $\alpha_3 = 1$  and  $\alpha_4 = 0.5$  for the synthetic tests and real events later. We also found that the final solutions are not very sensitive to small changes in the weights. The results for the synthetic test are summarized in Table 1. The first 200

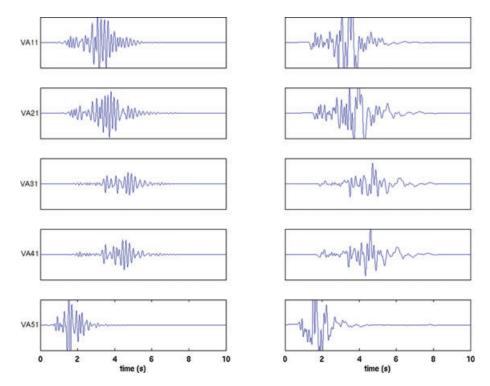


Figure 2. Synthetic seismograms (vertical component) at the five stations. The left column shows the filtered seismograms ( $3 \sim 9$  Hz), and the right column shows the unfiltered ones. The source is at 1227 m in depth, and has a strike of  $210^{\circ}$ , dip of  $50^{\circ}$  and rake of  $-40^{\circ}$ .

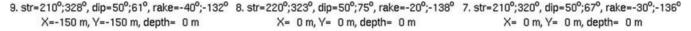
best solutions are used for statistical analysis of strikes, dips, rakes and locations. It is shown that even for the perfect data, the ambiguities (one standard deviation) in strike, dip and rake are about  $10^{\circ}$  because only vertical components are involved, and the station coverage is sparse. Among strike, dip and rake, dip has the least standard deviation while strike has the largest. As a small variation in depth changes the reflection and refraction patterns considerably, our algorithm determines the depth (*Z*) without much ambiguity.

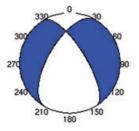
We further test the robustness of the method by adding noise to the synthetic data. We add white spectrum Gaussian noise to each trace, with zero mean and a standard deviation of 5 per cent of the maximum absolute amplitude of that trace. This level corresponds to the typical noise level we encounter for real data.

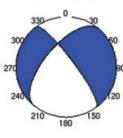
Fig. 3 shows the focal mechanism determined using waveform information and only three first P arrival polarities (we assume two polarities out of five are not identifiable due to noise contamination). The best solution here (#1) matches the correct solution. Fig. 4 shows the comparison between the modelled and synthesized waveforms with noise contamination. The 'shift' in the title of each subplot indicates the time shifted in the data to align with the synthetic waveforms. The reasons for having some time-shift

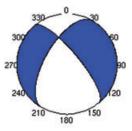
are as follows: (1) we introduced some artificial error in arrival time by manually picking the first P arrival in the synthetic data; (2) scattering noise can change the maximum cross-correlation position (Nolet *et al.* 2005). In the left column, the '+' or '-' signs indicate the first arrival polarities of P waves in the data and those in the synthetics; the upper ones are signs for the synthetic data while the lower ones are signs for the modelled data. The modelled traces all have the identical polarities as their counterparts in the synthetic data. Note that for the evaluation of the polarities, we use the unfiltered waveforms, as filtering usually blurs or distorts the polarities. In the right column, the number to the left of the slash denotes the S/P ratio for the modelled waveform. They are quite close in most cases.

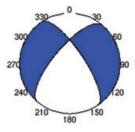
We further analysed the distribution of strikes, dips, rakes and locations from the first 200 best determined solutions (Table 2). The strike, dip and rake all have a mean quite close to the correct solution  $(210^\circ, 50^\circ \text{ and } -40^\circ, \text{ respectively})$ . Among the strike, dip and rake, dip has the least standard deviation. We find that in the *X* and *Y* directions the variation is much larger than that in the *Z* direction, similar to the case of clean data. The different standard deviations

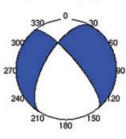


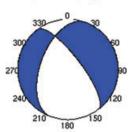




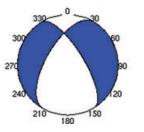


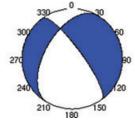






3. str=210°;328°, dip=50°;61°, rake=-40°;-132° 2. str=210°;324°, dip=40°;71°, rake=-30°;-126° 1. str=210°;328°, dip=50°;61°, rake=-40°;-132° X= 0 m, Y=-150 m, depth= 0 m X= 0 m, Y= 0 m, depth= 0 m X= 0 m, Y= 0 m, depth= 0 m





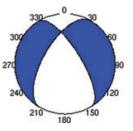


Figure 3. Nine best solutions from contaminated synthetic data; the number before 'str' is the order; '1' means the best solution and '9' means the worst solution among these nine.

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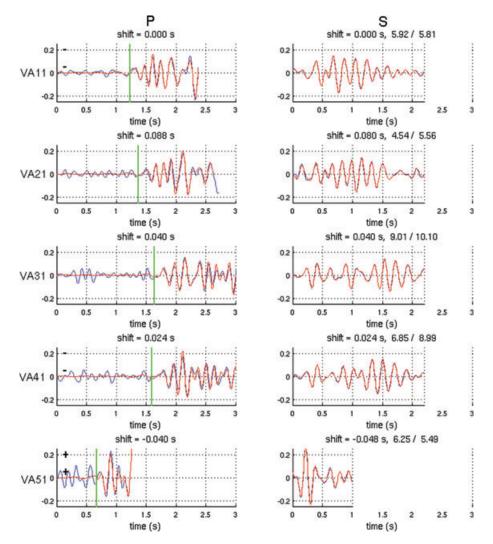


Figure 4. Comparison between modelled waveforms (red) and noisy synthetic data (blue) at five stations. From top to bottom waveforms from the vertical components at stations 1 to 5, respectively, are shown. The left column shows P waves and right column shows S waves. The green lines indicate the first P arrival times. For P waves, zero time means the origin time, and for S waves, zero time means the S-wave arrival time predicted by the calculated traveltime.

**Table 2.** Statistics of focal mechanism parameters in the 'noisy' synthetic test. The source location and mechanism are the same as the test in Table 1. The best solution has the identical focal parameters with the source.

	Strike (°)	Dip (°)	Rake (°)	<i>X</i> (m)	<i>Y</i> (m)	<i>Z</i> (m)
Mean	203	45	-41	-20	-10	0
Std.	13	10	11	120	130	0

might be an indication of the sensitivity of these model parameters. Therefore, using waveform information makes it possible to obtain accurate focal depth. We also tested other focal mechanisms with the same station coverage, and the recovered solutions are very close to the real ones, indicating that the method presented here is not sensitive only to certain event types due to the use of full waveform information.

### 4 APPLICATION TO INDUCED SEISMICITY AT AN OIL FIELD

We applied this method to study earthquakes at a petroleum field, which are induced by stress change due to water injection and gas/oil extraction. The earthquakes are small ( $-0.50 < M_w < 1.0$ ),

and the dominant energy in the recorded seismograms is between 3 and 15 Hz. Fig. 5 shows a typical event recorded at these stations and its spectrograms. During the period of 1999–2007, over 1500 induced earthquakes were recorded by a five-station near-surface network, and their occurrence frequency was found to be correlated with the amount of gas production (Sarkar 2008). Two major fault systems have been identified in this area, with one oriented in the NE–SW direction, and the other oriented in the conjugate direction (NW–SE) (Fig. 6).

The distribution of induced events in the field is shown in Fig. 6 (Sarkar 2008; Sarkar *et al.* 2008; Zhang *et al.* 2009). All the events have a residual traveltime less than 30 ms, indicating they are well located. First we describe the analysis procedure for one event. Fig. 7 shows the change of objective function value with the best solution order. For this event, the value is about 13 for the best solution and decreases to about 6 for the 200-th best solution. The objective function value decreases quickly from the 1st to the 200th best solution but relatively slowly beyond this range. Therefore, we choose 200 as the pool size for evaluating the statistics of focal mechanism parameters. The synthetic tests shown in the previous section have a similar objective function value distribution to the real data. Fig. 8 shows the beachballs of the nine best solutions

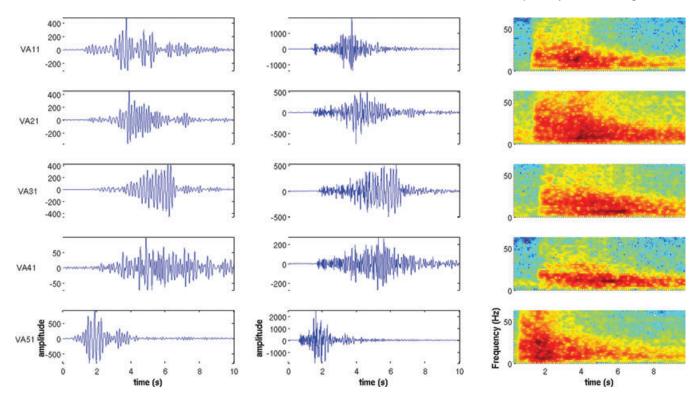


Figure 5. A typical event used in the focal mechanism determination and its spectrograms. The seismograms are from the vertical components of these five stations. The filtered seismograms ( $3 \sim 9$  Hz) are at the left column; the original seismograms are in the middle; the spectrograms of the original seismograms are at the right.

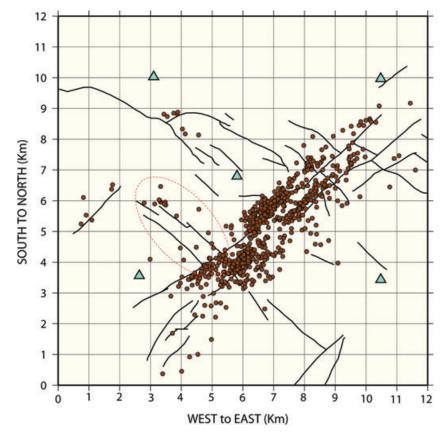


Figure 6. Distribution of over 1000 located induced earthquakes. Note that the majority of these earthquakes are in the proximity of the NE–SW fault. A group of induced earthquakes within the dashed circle may indicate the activation of a conjugate fault (Sarkar 2008). The five stations are indicated with green triangles, and station names are shown in Fig. 1(a).

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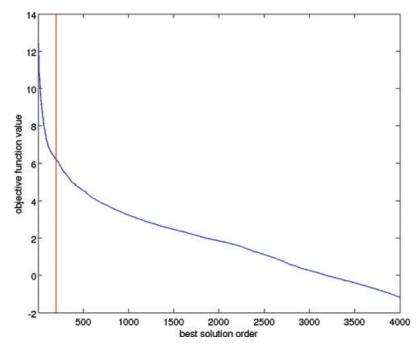
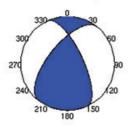
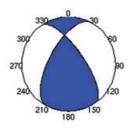


Figure 7. Objective function value versus best solution order. To the left of the red line are those solutions used to evaluate the statistics of focal mechanism and location parameters.

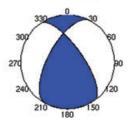
9. str=335°;218°, dip=60°;52°, rake=45°;141° 8. str=325°;199°, dip=60°;45°, rake=55°;135° 7. str=315°;198°, dip=60°;52°, rake=45°;141° X=150 m, Y=600 m, depth= 50 m X=150 m, Y=900 m, depth= 50 m X=300 m, Y=900 m, depth= 50 m

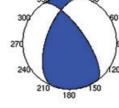


6. str=325°;208°, dip=60°;52°, rake=45°;141° X=150 m, Y=900 m, depth= 50 m



3. str=325°;208°, dip=60°;52°, rake=45°;141° 2. str=325°;208°, dip=60°;52°, rake=45°;141° X=300 m, Y=750 m, depth= 50 m





5. str=325°;188°, dip=60°;38°, rake=65°;126°

210

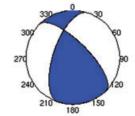
X=300 m, Y=750 m, depth= 50 m

180

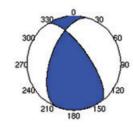
X=150 m, Y=750 m, depth= 50 m

150

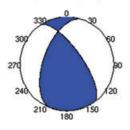
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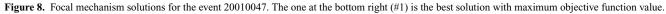


4. str=325°;199°, dip=60°;45°, rake=55°;135° X=150 m, Y=750 m, depth= 50 m



1. str=325°;199°, dip=60°;45°, rake=55°;135° X=300 m, Y=750 m, depth= 50 m





180

210

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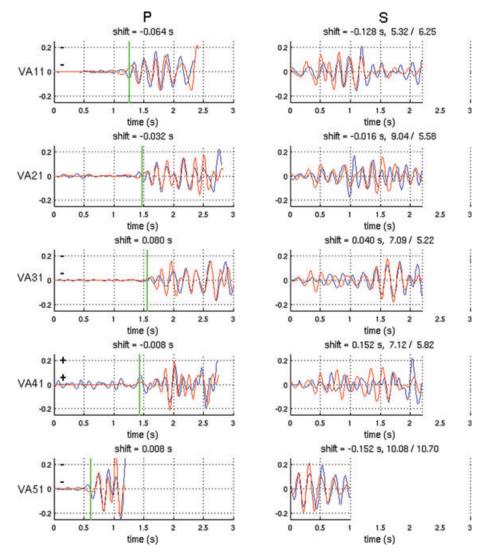


Figure 9. Comparison between the modelled waveforms (red) and the real data (blue) at five stations for the event 20010047. For *P* waves, zero time means the origin time, and for *S* waves, zero time means the *S*-wave arrival time predicted by the calculated traveltime.

**Table 3.** Statistics of focal mechanism parameters for the event 20010047. The best solution has a strike of  $325^{\circ}$ , dip of  $60^{\circ}$  and rake of  $55^{\circ}$ . In discussing the mean and standard deviation, the hypocentre is recoordinated as (0, 0, 0) m. The epicentre of event 20010047 is shown in Fig. 10 in a local coordinate.

	Strike (°)	Dip (°)	Rake (°)	X (m)	<i>Y</i> (m)	<i>Z</i> (m)
Mean	331	56	59	210	700	50
Std.	12	8	13	100	160	0

out of millions of trials. Our best solution (the one at the bottom right, reverse strike-slip) has a strike of 199°, which is quite close to the best known orientation 219° of the NE–SW fault (Fig. 1). However, the faulting (auxiliary direction  $325^{\circ}$ ) could occur in the conjugate NW–SE fault instead, as the auxiliary plane has a strike almost parallel with the conjugate fault. Using our new algorithm, the epicentre is shifted northwards by about 750 m, eastwards by about 300 m and the depth is shifted 50 m deeper. The shift in epicentre may be biased by errors in first *P* arrival picking or biased by inaccuracy in the velocity model. As has been discussed before,

the shift in epicentre may compensate the phase shift in the modelled seismograms due to inaccuracy in the velocity model.

Fig. 9 shows the comparison between the modelled and the observed data for event 20010047. The waveform similarities between the modelled and observed data are good. Additionally, the S/Pwaveform amplitude ratios in the modelled and observed data are quite close, and the first P arrival polarities are identical in the modelled and observed data for each station. In this example, all four criteria in eq. (2) are evaluated, and they are consistent between the modelled and observed data. In some situations, the observed and modelled waveforms may look similar, even when the S/P amplitude ratio and first P-wave polarity do not match. When this happens, our method would not accept this solution. This indicates that it is sometimes misleading to use only waveform matching to determine the focal mechanism, as a wrong solution might still give satisfactory waveform matching, especially when data from a sparse network are used. Table 3 shows the distribution of strike, dip, rake and location. Again, we find dip has the minimum standard deviation. Similar to the synthetic test, depth has the least variation in this case. As discussed before, variation in epicentre (X and Y) can be compensated by shifting the observed waveforms

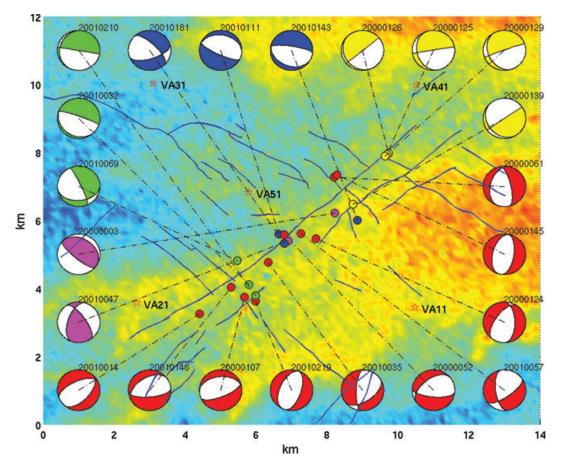


Figure 10. Focal mechanisms of the 22 events inverted in this study for an oil field. The colour in the map indicates the local change in surface elevation with a maximum difference of about 10 m. Different focal mechanisms are grouped in several colours.

to find a better alignment with the modelled data. Therefore, the constraint on lateral shift is weaker compared to that in the vertical direction.

Using this method, we have studied 22 earthquakes distributed along the NE–SW fault at this petroleum field. These 22 events are from the data collected between 2000 and 2002. The distribution of located induced earthquakes is shown in Fig. 6. The determined focal mechanisms are shown in Fig. 10. In this petroleum field, we expect some lateral velocity variations, especially across the faults (Zhang *et al.* 2009). We tested the robustness of our method by perturbing the *P* and *S* velocity model for each station independently to simulate lateral velocity heterogeneity (Appendix). The test showed that we are able to obtain a focal mechanism that is quite close to the correct solution when lateral velocity variation exists. This is because the new method incorporates different aspects of waveform information and compensates for the phase shift due to velocity heterogeneity using the time-shift.

Fig. 10 shows that the majority of the events primarily have the normal faulting mechanism, although there are also a few reverse faulting events. The dominance of the normal faulting mechanism suggests that the vertical stress is greater than the horizontal stress oriented in the NE–SW direction, parallel with strike of the main fault. Among these events, over half have a strike oriented in the NE–SW direction. However, the strike of some faults is in the conjugate fault direction (NW–SE). The dual orientations suggest that both fault systems are probably still active, and the seismicity pattern shown in Fig. 6 also supports this indication (e.g. the induced earthquakes in the red circle).

#### **5** CONCLUSIONS

In this study, we showed that combining the high-frequency seismograms with a fast optimized grid-search algorithm leads to determination of focal mechanisms and locations of small earthquakes, where subsurface velocity information is available. This method is especially applicable to the study of induced earthquakes recorded by a small number of stations, even when some first P arrival polarities are not identifiable due to noise contamination, or only the vertical components are usable. In addition, because of the normalization in our algorithm, the method does not require the absolute amplitude to be available and, therefore, mitigates the influence of the site effects. The objective function, formulated to include matching phase and amplitude information, first arrival P polarities and S/P amplitude ratios between the modelled and observed waveforms, yields stable solutions. The synthetic tests prove that the method is robust: giving correct solutions in the case of noise contamination or lateral velocity variations.

For real events, we find that the focal mechanisms are consistent with local geological structure and are indicative of local stress distribution. Focal mechanisms for 22 induced earthquakes are mostly normal faulting. A majority of the events are on the NE–SW striking faults, and their mechanisms are consistent with these faults. A few events have strikes in the NW–SE direction. These are events on the conjugate faults. In the region where both the NE trending faults and the conjugate (NW trending) faults exist, the focal mechanisms make it possible to determine with which faults the seismic events were associated.

#### ACKNOWLEDGMENTS

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#### APPENDIX: THE EFFECT OF THE VELOCITY MODEL ON THE FOCAL MECHANISM DETERMINATION

We test our new algorithm in the case when the 1-D velocity model is not a satisfactory approximation of the realistic subsurface structures, as lateral velocity variation exists locally. In general, we have a reliable velocity model for both P and S waves from well logs (Sarkar 2008). However, because lateral velocity heterogeneity is inevitable, the influence of inaccuracy in the velocity model on the focal mechanism determination also needs to be examined. We still use a 1-D layered model to generate the synthetic data from the event to a station. However, for each layer, a perturbation on the velocity equal to 5 per cent of the layer's reference velocity is added, and the perturbation is independent for all five stations. The density here is not perturbed in this test, as the velocity perturbation is dominant in determining the characteristics of the waveforms. Also, the layer thickness is not perturbed, as perturbation in either layer velocity or thickness generates equivalent phase distortions from each layer. The perturbations in both P and S velocity models are shown in Fig. A1. Note the perturbations in P and S velocity are independent of each other, and are also independent for each layer and station.

Fig. A1 shows that the variation in velocity for a certain layer from one station to another can be rather large. Considering the region in our study is less than 10 km × 10 km, and it is mainly composed of flat sedimentary layers, this variation should be a reasonable approximation of the real maximum lateral heterogeneity. Here we still use only three first *P* arrival polarities. Fig. A2 shows the best nine beachball solutions. Encouragingly, the solutions do not differ significantly from the correct one (strike =  $210^{\circ}$ , dip =  $50^{\circ}$  and rake =  $-40^{\circ}$ ). The strikes, dips and rakes in these solutions differ from the correct one by about  $10^{\circ}$  and one grid spacing (150 m) in the horizontal direction. All of these solutions find the correct depth.

Fig. A3 shows the comparison between the synthetic waveforms generated from the reference model and the synthetic observed

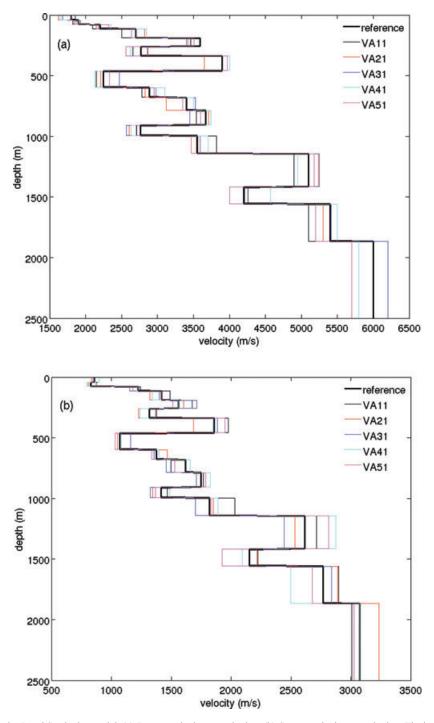
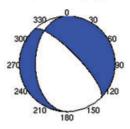


Figure A1. Perturbations in the *P* and *S* velocity model. (a) *P*-wave velocity perturbation. (b) *S*-wave velocity perturbation. The bold black line is the reference model, while the coloured lines are perturbed velocity structure for modelling seismograms in each station.

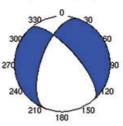
waveforms generated from the perturbed velocity model. Due to the different phase shift by velocity variation, many phases in the waveform have been distorted. However, allowing time-shift in the observed data compensates for much of the phase shift and distortion caused by velocity variation (note the 'shift' here is larger than in the previous case, especially for the *S*-wave comparison). Therefore, by incorporating information from different aspects in the waveform and compensating for the phase shift, we are still able to obtain a focal mechanism that is quite close to the correct solution. Table A1 shows the distribution of strike, dip and rake of the focal plane solution and X, Y and depth of the location from the first 200 best solutions. The mean values are rather close to the correct solutions, and the standard deviations are small, considering that only the vertical components of five stations and three polarities are used in determining the focal mechanism. Again, the depth variation is strongly constrained, and this indicates that using waveform information can greatly help locate the depth of an event, which is most difficult to constrain in the traditional traveltime-based location method. 9. str=200°;318°, dip=50°;61°, rake=-40°;-132° 8. str=200°;318°, dip=50°;61°, rake=-40°;-132° 7. str=180°;316°, dip=30°;67°, rake=-50°;-110° X=300 m, Y=150 m, depth= 0 m X=150 m, Y=150 m, depth= 0 m X=150 m, Y=150 m, depth= 0 m

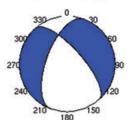


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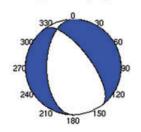




6. str=180°;323°, dip=40°;56°, rake=-60°;-113° 5. str=170°;316°, dip=30°;64°, rake=-60°;-106° 4. str=180°;323°, dip=40°;56°, rake=-60°;-113° X=150 m, Y=150 m, depth= 0 m X=150 m, Y=150 m, depth= 0 m X=150 m, Y= 0 m, depth= 0 m

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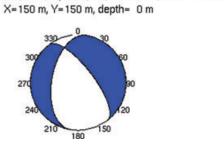
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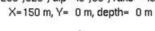
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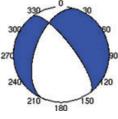
3. str=190°;322°, dip=40°;61°, rake=-50°;-118° 2. str=200°;323°, dip=40°;66°, rake=-40°;-123° 1. str=190°;322°, dip=40°;61°, rake=-50°;-118°

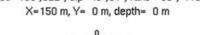
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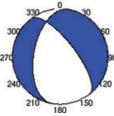


Figure A2. Focal mechanism solutions when lateral velocity variation exists. The best solution here (#1) does not perfectly match the correct solution but is quite close.

Table A1. Statistics of focal mechanism parameters in the velocity perturbed synthetic test. The source location and focal mechanism are the same as the clean synthetic test.

	Strike (°)	Dip (°)	Rake (°)	<i>X</i> (m)	<i>Y</i> (m)	<i>Z</i> (m)
Mean	194	42	-46	80	10	0
Std.	17	10	14	150	170	0

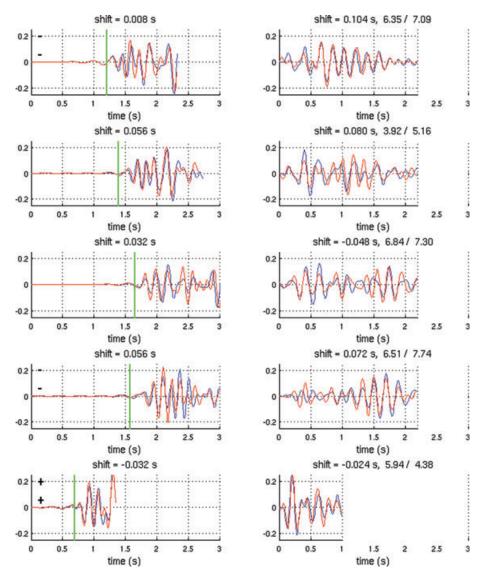


Figure A3. Comparison between the modelled and synthetic data when lateral velocity variation exists. Phase mismatch between the modelled and synthetic data is similar to the difference in the real event 20010047.